

# Dark Matter Bodies In The Solar Structure

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1. Data - solar surface oscillations observed at the frequencies

$$\nu_{CrAO} = 104.1891 \mu\text{Hz} \text{ (160 minutes)}$$

$$\nu_{SoHO} = 220.7 \mu\text{Hz} \text{ (75 minutes)}$$

$$V_{CrAO} = 0.25 \pm 0.03 \text{ m/s}$$

$$V_{SoHO} = 0.0045 \pm 0.0015 \text{ m/s}$$

CrAO (Crimean Astrophysical Observatory),

SoHO (Solar Helioseismic Observatory)

2. Astroseismology and resonant excitations of Surface Rayleigh Waves by two Dark Matter Bodies (DMB) with masses  $M_{CrAO} \approx 1.2 \cdot 10^7 M_{\oplus}$  ( $\approx 0.73 M_{\oplus}$ ),  $M_{SoHO} \approx 3.7 \cdot 10^9 M_{\oplus}$  ( $\approx 0.02 M_{\oplus}$ ) and radii of matter distributions

$$R_{CrAO} \approx 20 - 200 \text{ km},$$

$$R_{SoHO} \approx 60 - 400 \text{ km} \text{ orbiting near the solar surface at radii}$$

$$R_{OrbCrAO} = 1.54 R_{\oplus},$$

$$R_{OrbSoHO} = 0.94 R_{\oplus}.$$

3. These DMB radiate Gravitational Waves (GW) with amplitudes of metric tensor oscillations large enough to detect

$$h_{CrAO} \sim 9.5 \cdot 10^{-20}, \quad h_{SoHO} \sim 1.2 \cdot 10^{-21} \text{ near the Earth (near-zone) at frequencies } \nu_{CrAO}, \nu_{SoHO} \text{ and } \underline{\text{could be detected by}}$$

eLISA (European Laser Interferometer Space Antenna) or by

DULKYN-3 (Scientific Centre for Gravitational-wave Research "Dulkyn") at observation time much longer than

$$8 \text{ days at } \nu_{CrAO}, \quad 5 - 10 \text{ years at } \nu_{SoHO}, \quad 9 \text{ years for PSR J0737-3039 } (h_{PSR J0737-3039} \sim 2 \cdot 10^{-22}).$$

4. Suggested experiment – a possible solution of two fundamental physical problems simultaneously –

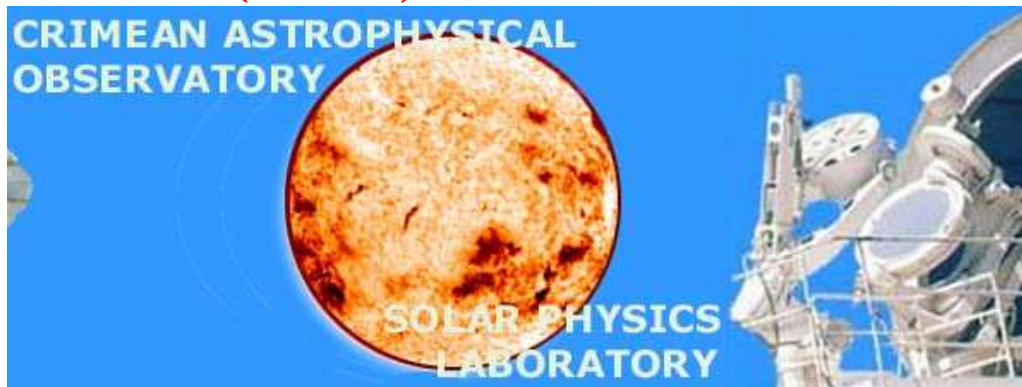
the direct detection of GW, and the direct detection of DMB in the Sun.

The first theoretical estimations – “Excitation of the solar oscillations by objects consisting of y-matter”

by S. I. Blinnikov, M. Yu. Khlopov in Solar Physics 82(1983)383

## Current Detectors:

### CrAO (1974)



### SoHO (1996)

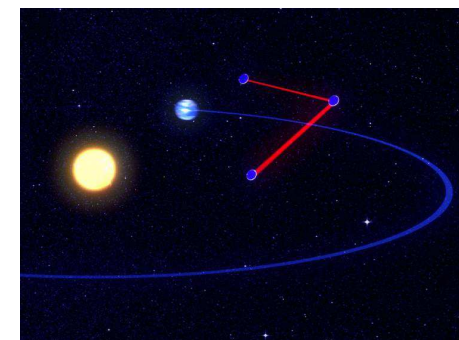
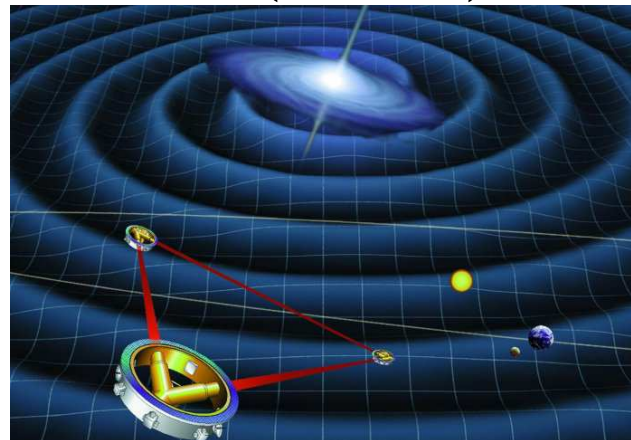


## Future Detectors:

### DULKYN-3 (2015?) ~ 1 m

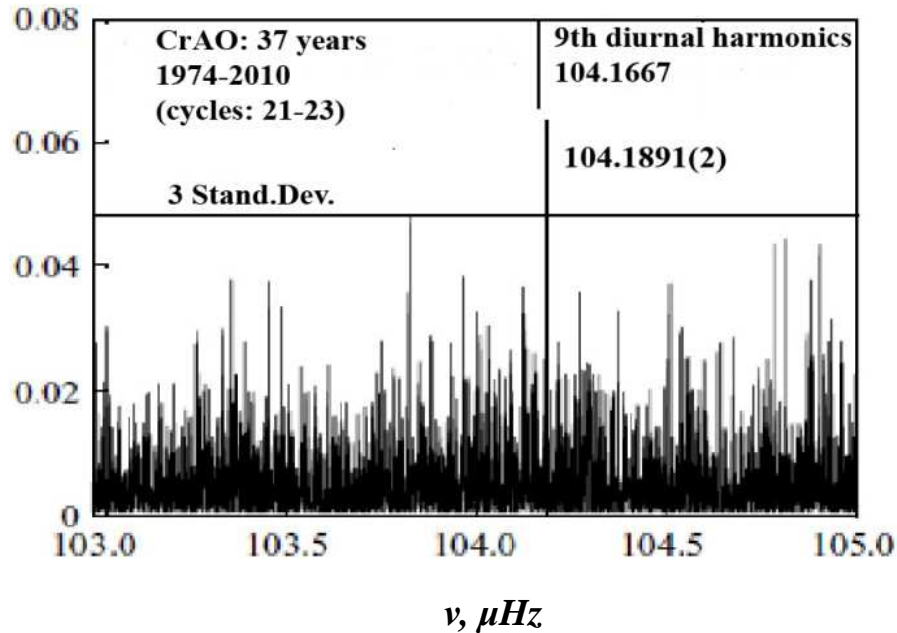


### eLISA (2016?) ~ 1000000 km



# Solar oscillation data

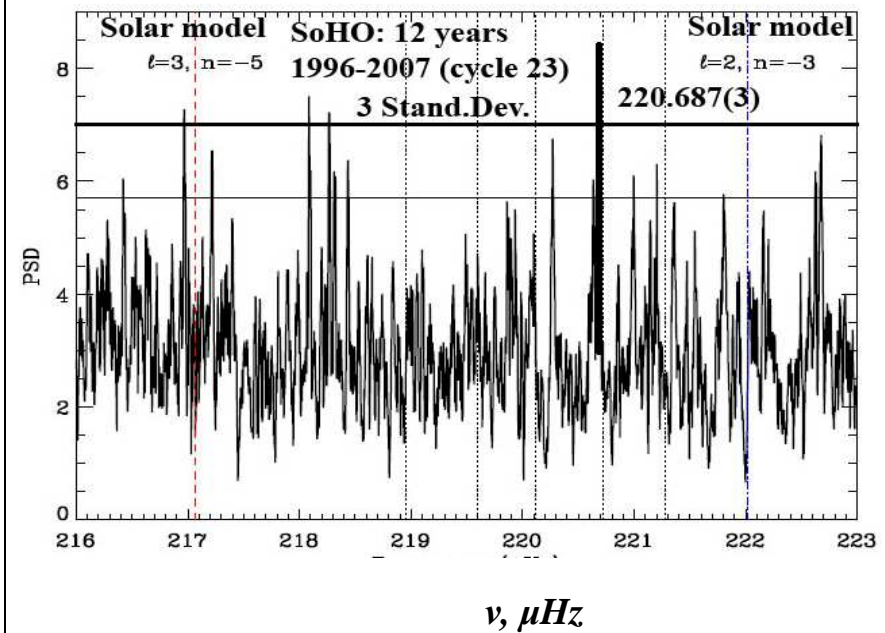
## CrAO (Crimean Astrophysical Observatory)



Power of the Fourier spectrum density extracted from differential (center - limb) velocity spectrum of solar surface oscillations (CrAO data) in 1974–2010 (21, 22, 23 solar cycles).

“Oscillations of the Sun: Results of observations in 1974–2007” by V. A. Kotov, V. I. Haneychuk, Bulletin of Crimean Astrophysical Observatory, 104 (2008) 45; “Pulsations of the Sun and a Beat Period of 399 Days” by V. A. Kotov, V. I. Khaneichuk, Bulletin of the Crimean Astrophysical Observatory 107 (2011) 67

## SoHO (Solar Helioseismic Observatory)



Power of the Fourier spectrum density extracted from intensity spectrum of solar surface oscillations (VIRGO/SPM at SoHO - blue channel data) in 1996-2007 (23 solar cycle). The vertical dashed lines at  $\approx 217$  and  $222 \mu\text{Hz}$  are the expected theoretical frequencies for the g-modes  $(\ell, n) = (2, -5)$  and  $(3, -2)$  from the Saclay seismic model.

“Update on g-mode research” by R.A. Garcia, A. Jimenez, S. Mathur, J. Ballot, A. Eff-Darwich, S.J. Jimenez-Reyes, P.L. Palle, J. Provost, S. Turck-Chieze. Astron. Nachr. AN 999(2006)88; arXiv:0802.4296 [astro-ph]

The first theoretical estimations – “Excitation of the solar oscillations by objects consisting of  $\gamma$ -matter” by S.I. Blinnikov, M.Yu. Khlopov, Solar Physics 82(1983)383

Equations of Astroseismology with tidal excitation by DMB +

Surface Boundary Conditions: Lagrange variations  $\delta p'(R) = 0$ ,  $\delta(\delta p'(R)) = 0 \Rightarrow$  Resonances like Surface Rayleigh Waves.

G-mode Resonances? - Do not appear at the star surface. For DMB orbit  $\mathbf{R}_{orb}(t) = \left\{ R_{orb}(t), \frac{\pi}{2}, \Phi(t) \right\}$ ,

$$U(r, t) = - \frac{G m_{DMB}}{|r - R_{orb}(t)|} = \frac{G m_{DMB}}{r_{>}} \sum_{l,lz} W_{l,lz} \left( \frac{r_{<}}{r_{>}} \right)^l e^{-i l z \Phi(t)} Y_l^{lz}(\theta, \varphi)$$

$$W_{l,lz} = Re \left( (-1)^{\frac{l+lz}{2}} \frac{\left[ \frac{4\pi}{2l+1} (l+lz)! (l-lz)! \right]^{1/2}}{\left[ 2^l \left( \frac{l+lz}{2} \right)! \left( \frac{l-lz}{2} \right)! \right]} \right), \quad W_{2,2} = \sqrt{\frac{3\pi}{10}}, \quad W_{6,\{2,4,6\}} = \left\{ \sqrt{\frac{105\pi}{3328}}, -\sqrt{\frac{126\pi}{3328}}, \sqrt{\frac{231\pi}{3328}} \right\}$$

$$\xi(r, t) = \sqrt{4\pi} \sum_{n,l,lz} C_{n,l,lz} \left\{ \xi_r(r) Y_l^{lz}(\theta, \varphi), \frac{1}{\sqrt{l(l+1)}} \xi_h(r) \frac{d}{dr} Y_l^{lz}(\theta, \varphi), \frac{1}{\sqrt{l(l+1) \sin \theta}} \xi_h(r) \frac{d}{d\varphi} Y_l^{lz}(\theta, \varphi) \right\} e^{-i \omega_{n,l,lz} t},$$

$$\xi_r(R) = - \frac{1}{\omega_{Dyn}^2} \frac{\frac{1}{R} U_{l,lz}(R) + \frac{1}{l(l+1)} \frac{\omega^2}{\omega_{Dyn}^2} U'_{l,lz}(R)}{1 - \frac{4}{(l+1)(l+2)} \frac{\omega^2}{\omega_{Dyn}^2} \left( 1 + \frac{1}{4} \frac{\omega^2}{\omega_{Dyn}^2} - \frac{\pi G \rho(R)}{\omega_{Dyn}^2} + \frac{1}{l+2} \frac{\pi G R \rho'(R)}{\omega_{Dyn}^2} \right)},$$

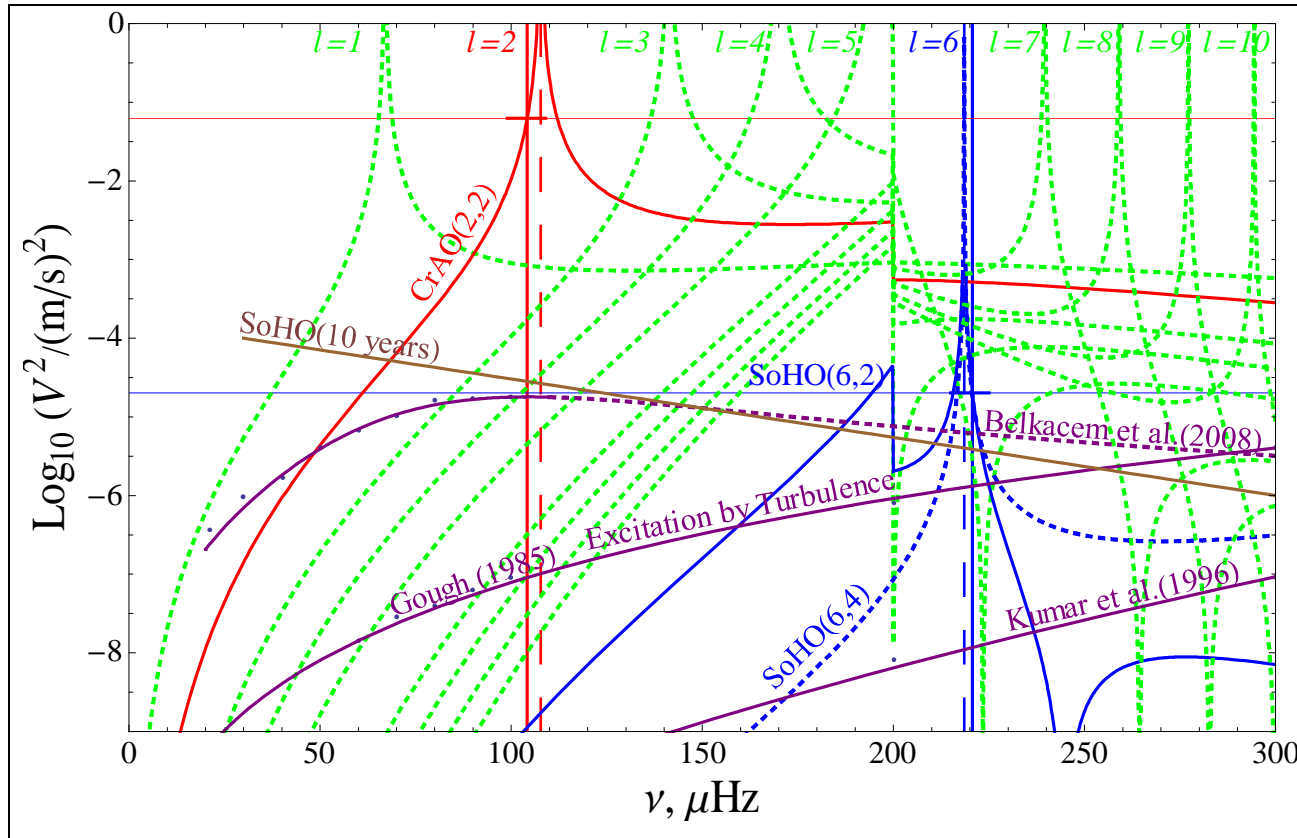
$$\xi_h(R) =$$

$$- \frac{\frac{4}{l(l+1)} \frac{1}{\omega_{Dyn}^2} \frac{1}{R} \left( 1 + \frac{1}{4} \frac{\omega^2}{\omega_{Dyn}^2} - \frac{\pi G \rho(R)}{\omega_{Dyn}^2} + \frac{1}{l+2} \frac{\pi G R \rho'(R)}{\omega_{Dyn}^2} \right) U_{l,lz}(R) + \frac{1}{4} \left( 1 + \frac{4}{(l+1)(l+2)} \frac{\pi G R \rho'(R)}{\omega_{Dyn}^2} \right) U'(R)}{1 - \frac{4}{(l+1)(l+2)} \frac{\omega^2}{\omega_{Dyn}^2} \left( 1 + \frac{1}{4} \frac{\omega^2}{\omega_{Dyn}^2} - \frac{\pi G \rho(R)}{\omega_{Dyn}^2} + \frac{1}{l+2} \frac{\pi G R \rho'(R)}{\omega_{Dyn}^2} \right)},$$

$$\omega_{Dyn}^2 = \frac{GM}{R^3}, \quad U_{l,lz}(r) = \frac{G m_{DMB}}{r_{>}} \sum_{l,lz} W_{l,lz} \left( \frac{r_{<}}{r_{>}} \right)^l - \text{for a circular orbit. For solar like stars: } \frac{G \rho(R)}{\omega_{Dyn}^2} \ll 1, \frac{G R \rho'(R)}{\omega_{Dyn}^2} \ll 1,$$

$$\omega^2 \rightarrow \omega_{l,l_z}(\omega_{l,l_z} - l_z \Omega_{star}) \Rightarrow \omega_{l,l_z} = \omega_{Dyn} \sqrt{\sqrt{l(l+1)+4} - 2 + \frac{l_z^2 \Omega_{star}^2}{4 \omega_{Dyn}^2}} + \frac{1}{2} l_z \Omega_{star} \Rightarrow \omega_l = \omega_{Dyn} \sqrt{\sqrt{l(l+1)+4} - 2}$$

## Amplitudes of solar surface oscillations



**For the Sun:**

$$\frac{\Omega_{Sun}^2}{\omega_{Dyn}^2} \ll 1 \Rightarrow$$

$$\omega_{l,l_z} = \omega_{Dyn} \sqrt{\sqrt{l(l+1)+4} - 2 + \frac{1}{2} l_z \Omega_{Sun}}$$

$$v_{Dyn} = \frac{\omega_{Dyn}}{2\pi} = 99.9178 \mu\text{Hz} \quad (166.804')$$

$$\underline{M_{CrAO} \approx 1.2 \cdot 10^{-7} M_{\odot} (\approx 0.73 M_{\oplus}), l=2, l_z=2}$$

$$\underline{M_{SoHO} \approx 3.7 \cdot 10^{-9} M_{\odot} (\approx 0.02 M_{\oplus}), l=6, l_z=2,4,6.}$$

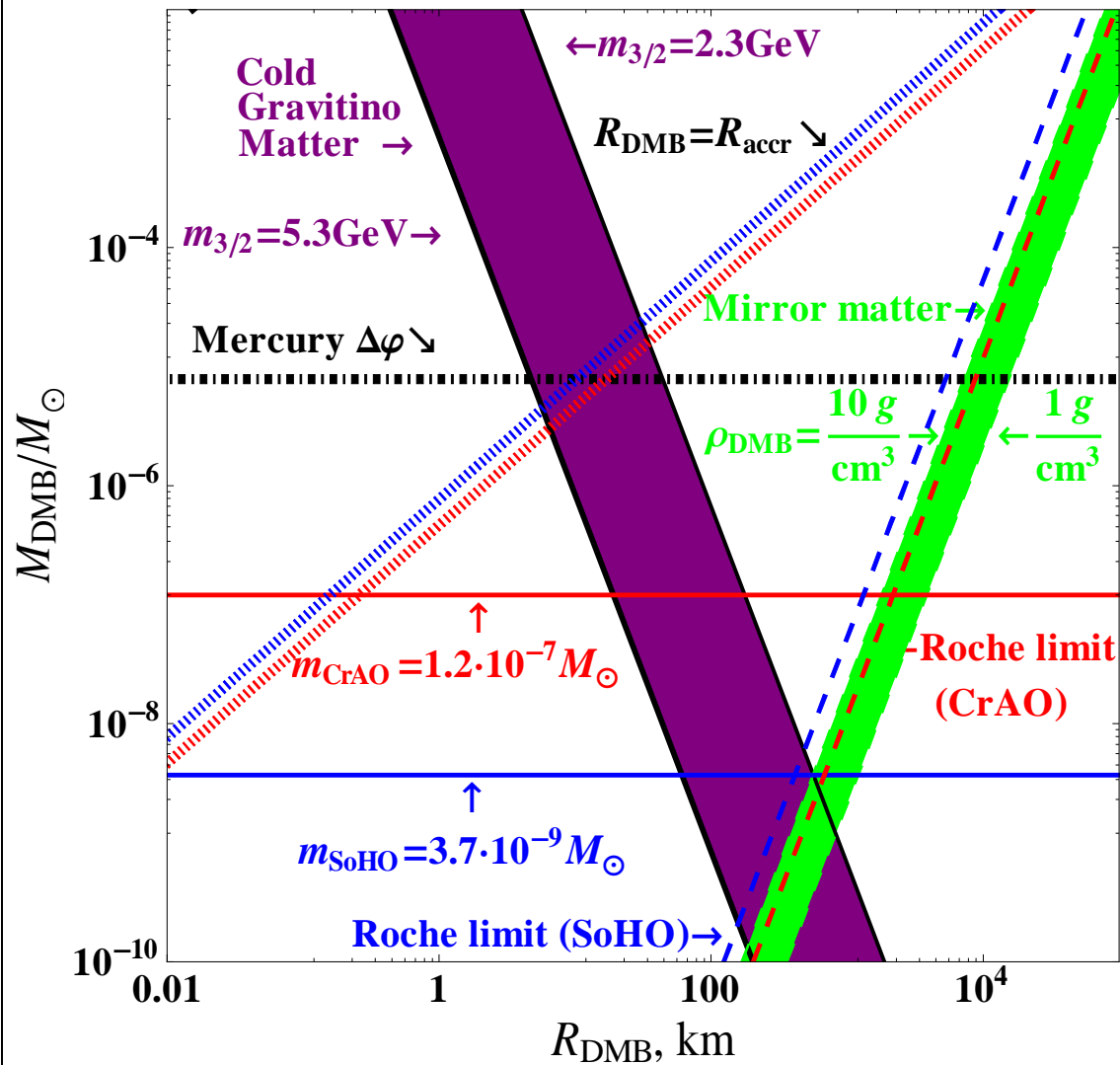
**More accurate solar surface experiments or Gravitational Wave experiments are needed**

D. O. Gough in E. Rolfe, B. Battrick (eds.), Future Missions in Solar, Heliospheric and Space Plasma Physics, ESA SP-235, ESA Publications Division, Noordwijk, The Netherlands, (1985) 183

P. Kumar, E. J. Quataert, J. N. Bahcall, "Observational searches for solar g modes: Some theoretical considerations", ApJ Letters 458 (1996) L83

K. Belkacem, R. Samadi, M. J. Goupil, M. A. Dupret, A. S. Brun, and F. Baudin, "Stochastic excitation of nonradial modes II. Are solar asymptotic gravity modes detectable?", Astronomy & Astrophysics manuscript no. 0827 c ESO 2008

**Mass – Radius Limitations for DMB**



$M_{\text{DMB}} \ll 8 \cdot 10^{-6} M_{\odot}$   
 Mercury Perihelion Precession

$R_{\text{DMB}} \ll R_{\text{RocheLimit}}$

$R_{\text{DMB}} \gg R_{\text{accr}}$

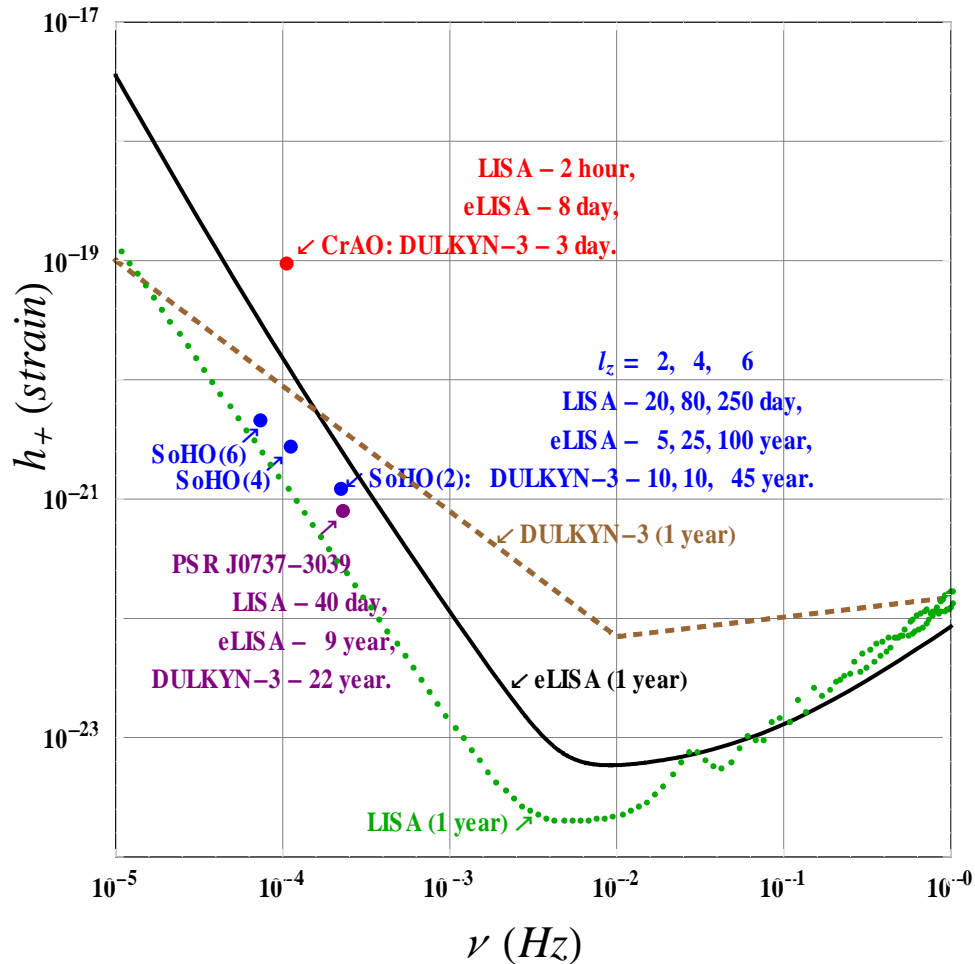
**Cold Gravitino Matter?**

## Estimations for future GW experiments

$$h_+ = -\frac{1}{R} \frac{G^2}{c^4} \frac{2m_1 m_2}{r} (1 + \cos^2 \theta) \cos [2\omega(t - R)],$$

$$h_\times = -\frac{1}{R} \frac{G^2}{c^4} \frac{4m_1 m_2}{r} (\cos \theta) \sin [2\omega(t - R)].$$

+ near zone terms



## Conclusions:

**CrAO** and **SoHO** data for oscillation of solar surface indicate two DMB in the solar structure at

Mass of DM particles  $\ll 30 \text{ GeV}/c^2$

$$\underline{M_{CrAO} \approx 1.2 \cdot 10^7 M_\odot (\approx 0.73 M_\oplus)} \quad \underline{M_{SoHO} \approx 3.7 \cdot 10^9 M_\odot (\approx 0.02 M_\oplus)}$$

$$\underline{R_{CrAO} \approx 20 - 200 \text{ km}},$$

$$\underline{R_{CSOHO} \approx 60 - 400 \text{ km}}$$

$$\underline{R_{OrbCrAO} = 1.54 R_\odot},$$

$$\underline{R_{OrbSoHO} = 0.94 R_\odot}.$$

At the Earth the Sun is the brightest source of GW at frequencies

$$\underline{v_{CrAO} = 104.1891 \mu\text{Hz}}$$

$$\underline{v_{SoHO} = 220.7 \mu\text{Hz}}$$

$$\underline{h_{CrAO} \sim 9.5 \cdot 10^{-20}}$$

$$\underline{h_{SoHO} \sim 1.2 \cdot 10^{-21}}$$

$$\underline{v_{PSR J0737-3039} = 231.5 \mu\text{Hz}}$$

$$\underline{h_{PSR J0737-3039} \sim 2 \cdot 10^{-22}}$$

**eLISA (2016?)**

**DULKYN-3 (2015?)**

8 days at  $v_{CrAO}$

3 days at  $v_{CrAO}$

5, 25, 100 years at  $v_{SoHO}$

10, 10, 45 years at  $v_{SoHO}$

9 years at  $v_{PSR J0737-3039}$

22 years at  $v_{PSR J0737-3039}$

GW experiments **to detect at least two DMB** in the Sun, **to fix the  $l_z = 2, 4, 6$  uncertainty** for SoHO oscillations, and **to detect other possible DMB in the Sun** which do not appear in surface oscillations because their frequencies are too far from the surface resonances and amplitudes are too small to detect.

**Suggested experiment** – a possible solution of two fundamental physical problems simultaneously – **the direct detection of GW**, and **the direct detection of DMB in the Sun**.

In the case of success, this experiment opens a **new epoch in Astro, Particle and GW physics**.

