

Top quark physics

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Why we like the t -quark ?

The top-quark is an extraordinary Standard Model (SM) object

- the most heavy SM object, $m_t \approx 172.5$ GeV and $|V_{tb}| \lesssim 1$
- \Rightarrow top decays before hadronization \Rightarrow there are no "top" hadrons
- connection with other generations is very small: $|V_{td}| \simeq 0.008$, $|V_{ts}| \simeq 0.04$
- extraordinary accuracy of the theoretical predictions ($\sigma, \Gamma, \text{Br}, \dots, \sim \mathcal{O}(1\%)$)
- all properties are described within the SM without additional phenomenological parameters
- provides the direct information about spin and polarization
- the largest Yukawa coupling: $y_t = \sqrt{2}m_t/v \approx 1$

t -quark is an excellent laboratory to search for New Physics

- new interactions and particles
- new anomalous interactions: tWb ; tHq ; $tg/\gamma/Zq$
- new heavy objects $R(t\bar{t})$, $Q \rightarrow tX, \dots$
- ...

Top-quark interactions within the Standard Model

The Lagrangian

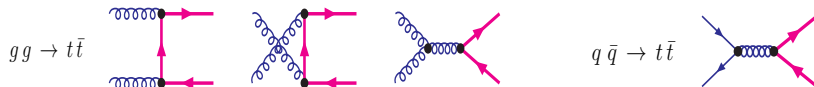
$$\begin{aligned} \mathcal{L}_{\text{SM}} = & -\frac{y_t}{\sqrt{2}} \bar{t} t H - g_s \bar{t} \gamma^\mu t^a t G_\mu^a - \frac{g}{\sqrt{2}} \sum_{q=d,s,b} \frac{V_{tq}}{2} \bar{t} \gamma^\mu (1 - \gamma_5) q W_\mu^+ \\ & - Q_t e \bar{t} \gamma^\mu t A_\mu - \frac{g}{2 \cos \vartheta_W} \bar{t} \gamma^\mu \left[\left(\frac{1}{2} - 2Q_t \sin^2 \vartheta_W \right) - \frac{1}{2} \gamma_5 \right] t Z_\mu + \text{h.c.} \end{aligned}$$

$$y_t = \sqrt{2} \frac{m_t}{v_{\text{ew}}} \approx 1, \quad v_{\text{ew}} \approx 246 \text{ GeV}$$

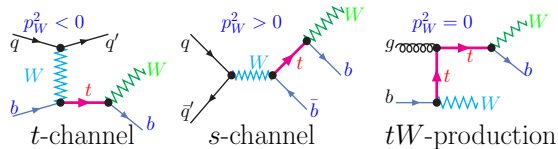
v_{ew} - electroweak scale - vacuum expectation value

Top-quark production processes in the SM

QCD t -quarks production



t -quarks production due to electroweak interactions



Top-quark decays within the SM

dominant decay channel $t \rightarrow b W^+; W^+ \rightarrow q \bar{q}', \ell \nu$

decay width (neglecting m_b^2/m_t^2)

$$\Gamma_{tot} = \Gamma_t = \frac{G_F m_t^3}{8\sqrt{2}\pi} \left(1 - \frac{M_W^2}{m_t^2}\right)^2 \left(1 + 2\frac{M_W^2}{m_t^2}\right) \left[1 - \frac{2\alpha_s}{3\pi} \left(\frac{2\pi^2}{3} - \frac{5}{2}\right)\right]$$

$$G_F = 1.16637 \times 10^{-5} \text{ GeV}^{-2}, \quad m_t = 172.5 \text{ GeV}, \quad \alpha_s(m_t) = 0.118$$

$$\Rightarrow \Gamma_t \simeq 1.39 \text{ GeV} \gg \Lambda_{QCD} \simeq 200 \text{ MeV}$$

t -quark decays before hadronization. There are no "top"-hadrons ($t\bar{t}$), ($t\bar{q}$), (tqq')

$$BR(t \rightarrow b \ell^+ (e, \mu, \tau) \nu) \simeq 33\%; \quad BR(t \rightarrow b q q') \simeq 67\%$$

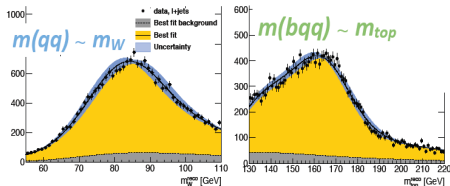
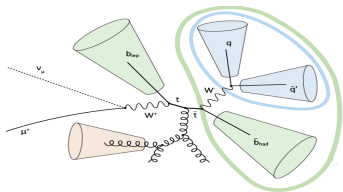
decay probabilities for $t\bar{t}$ pair

$$\begin{aligned} t\bar{t} &\rightarrow b \bar{b} q \bar{q}' q'' \bar{q}''' && \simeq 46.2\% \\ t\bar{t} &\rightarrow b \bar{b} q \bar{q}' \ell^+ \nu + b \bar{b} q \bar{q}' \ell^- \bar{\nu} && \simeq 43.5\% \\ t\bar{t} &\rightarrow b \bar{b} \ell^+ \ell^- \nu \bar{u} && \simeq 10.3\% \end{aligned}$$

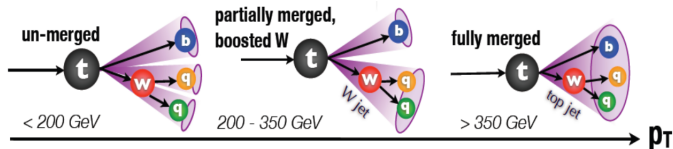
the final states have two b -jets; 0, 1, 2 "isolated" charged leptons; "missing" energy (neutrino) and several hadronic jets from light quarks

Top-quark reconstruction

- two ways for the reconstruction: $t \rightarrow bW^+(\rightarrow \ell^+\nu)$ and $t \rightarrow bW^+(\rightarrow q\bar{q}')$
reconstruction of the $W(m(jj) \sim M_W)$, then $m(J_b "W") = m(j_b jj)$



- "boosted" t -quark

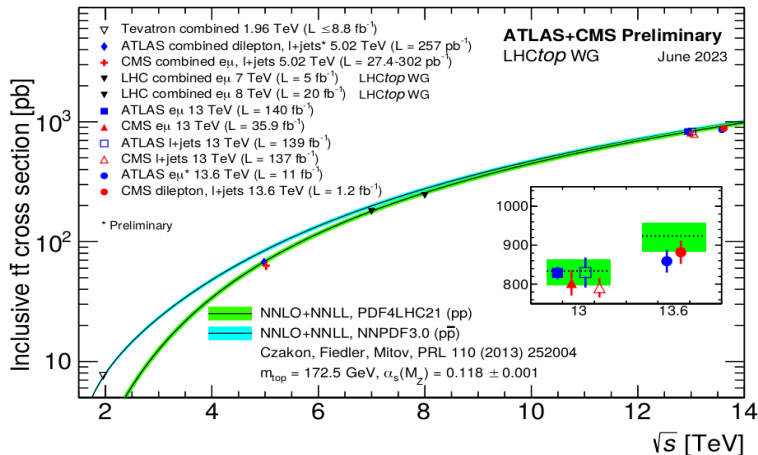


Total cross sections production of $t\bar{t}$ at the LHC

Great progress has been made in describing the production of $t\bar{t}$ taking into account higher orders of perturbation theory (NNLO and NNLL). The main theoretical uncertainty comes from PDF

The cross sections for the production of the pair $t\bar{t}$ are measured in various final states:

$ee, \mu\mu, e\mu, \ell\tau_h, e/\mu + \text{jets}, \text{all jets}$



Total cross sections production of $t\bar{t}$ at the LHC

joint analysis of two experiments ("LHC Top Physics Working Group"):

| \sqrt{s} , TeV | $\sigma(\text{pb})$ (experiment) | $\sigma(\text{pb})$ (theory) |
|------------------|--|--|
| 7 | $69.5 \pm 6.1(\text{stat}) \pm 5.6(\text{syst}) \pm 1.6(\text{lumi})$ | 70 ± 10 (scale + PDF + α_s) |
| 7 | $173 \pm 3(\text{stat}) \pm 8(\text{syst}) \pm 6(\text{lumi})$ | $177.3^{+4.7}_{-6.8}(\text{scale}) \pm 9$ (PDF + α_s) |
| 8 | $240.6 \pm 1.4(\text{stat}) \pm 5.7(\text{syst}) \pm 6.2(\text{lumi})$ | $252.9^{+6.4}_{-8.6}(\text{scale}) \pm 11.7$ (PDF + α_s) |
| 13.0 | $836 \pm 27(\text{stat}) \pm 81(\text{syst}) \pm 100(\text{lumi})$ | $832.0^{+20}_{-28}(\text{scale}) \pm 35$ (PDF + α_s) |

- there are two measurements of the top pair production in pp -collisions at $\sqrt{s} = 13.6$ TeV

CMS: JHEP 08, p. 204 (2023), arXiv:2303.10680, $\mathcal{L}_{int} = 1.21 \text{ fb}^{-1}$

ATLAS: arXiv:2308.09529, $\mathcal{L}_{int} = 29 \text{ fb}^{-1}$

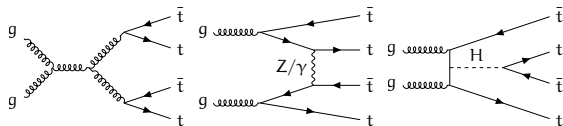
ATLAS $\sigma(t\bar{t}) = 850 \pm 3(\text{stat}) \pm 18(\text{syst}) \pm 20(\text{lumi}) \text{ pb}$

CMS $\sigma(t\bar{t}) = 881 \pm 23(\text{stat+syst}) \pm 20(\text{lumi}) \text{ pb}$

theory $\sigma(t\bar{t}) = 924^{+32}_{-40}(\text{scale} + \text{PDF} + \alpha_s) \text{ pb}$

Four-top-quark production

Four top quark production in pp -collisions is among the rarest SM processes currently accessible at hadron colliders



The SM cross section is calculated at NLO in QCD and EW theory *JHEP 02 (2018) 031*, *arXiv:1711.02116 [hep-ph]* and *arXiv:2212.03259 (2022)*

$$\sigma^{theory} = 12.0 \pm 2.4 \text{ fb and } = 13.4_{-1.8}^{+1.0} \text{ fb } \sqrt{s} = 13 \text{ TeV}$$

there two measurements of the four-top-quark production in pp -collisions at $\sqrt{s} = 13 \text{ TeV}$

ATLAS: *Eur. Phys. J. C 83 (2023) 496*, *arXiv:2303.15061*

$$\mathcal{L}_{int} = 140 \text{ fb}^{-1}, \quad \sigma(t\bar{t}t\bar{t}) = 22.5_{-4.3}^{+4.7} (\text{stat})_{-3.4}^{+4.6} (\text{syst}) \text{ pb} = 22.5_{-5.5}^{+6.6} \text{ fb}$$

top-quark Yukawa coupling $|\kappa_t| < 2.2(1.8)$

CMS: *Phys.Lett.B 847 (2023) 138290*, *arXiv:2305.13439 [hep-ex]*

$$\mathcal{L}_{int} = 138 \text{ fb}^{-1}, \quad \sigma(t\bar{t}t\bar{t}) = 17.7_{-3.5}^{+3.7} (\text{stat})_{-1.9}^{+2.3} (\text{syst}) \text{ pb} = 17.4_{-4.0}^{+4.4} \text{ fb}$$

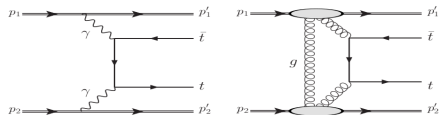
$$\sigma(t\bar{t}W) = 990 \pm 98 \text{ fb}, \quad \sigma^{theor} = 722 \pm 74 \text{ fb}$$

$$\sigma(t\bar{t}Z) = 945 \pm 81 \text{ fb}, \quad \sigma^{theor} = 859 \pm 80 \text{ fb}$$

top-quark Yukawa coupling $|\kappa_t| < 1.7$

$pp \rightarrow pt\bar{t}p$ and $PbPb \rightarrow t\bar{t}$

- CMS+TOTEM Collaborations, arXiv:2310.11231. Central exclusive production of $t\bar{t}$ pairs ...



fractional momentum loss of the intact protons $\zeta = (|\vec{p}_i| - |\vec{p}_0|)/|\vec{p}_i|$, \vec{p}_i and \vec{p}_0 are the momenta of the incoming and outgoing protons: $0.02 < \zeta < 0.20$

$$pp \rightarrow pt\bar{t}p : \sigma(t\bar{t}) < 0.59 \text{ pb at } 95\%, \quad \sqrt{s} = 13 \text{ TeV}$$

$$pp \rightarrow p\gamma\gamma p \rightarrow pt\bar{t}p : \sigma(t\bar{t})^{theor} = 0.22 \pm 0.05 \text{ fb}$$

- CMS Collaboration "Evidence for top quark production in nucleus-nucleus collisions", PRL 125 (2020) 22, 222001, arXiv:2006.11110

$$PbPb \rightarrow t\bar{t}X, \quad \sqrt{s_{NN}} = 5.02 \text{ TeV}$$

$$\sigma_{t\bar{t}} = 2.54_{-0.74}^{+0.84} (\ell\ell') = 2.03_{-0.64}^{+0.71} (+b) \mu\text{b}$$

it is compatible with previous CMS result times $A(Pb)^2$

$$pp \rightarrow t\bar{t}X, \quad \sqrt{s} = 5.02 \text{ TeV}, \quad \sigma_{t\bar{t}} = 69.5 \pm 8.43 \text{ pb}$$

Top-quark mass

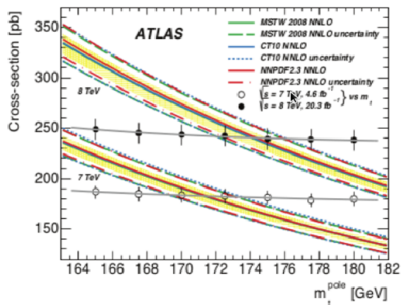
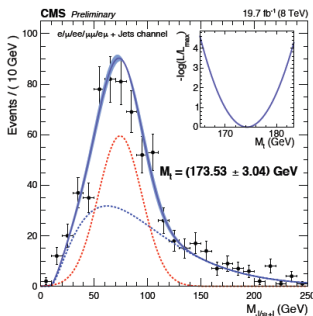
m_t is the fundamental parameter of theory

◇ related to other EW parameters – stringent tests of SM

◇ vacuum stability depends on exact value of m_t

◇ important for m_W , $\text{Br}(B_s \rightarrow \mu\mu)$, ...

- $m_t \leftarrow M_{inv}(j_b'' W''(jj))$
- $m_t \leftarrow M_{inv}(j_b M_T(\ell E_T^{miss}))$
- $m_t \leftarrow M_{inv}(j_b \ell^\pm), M_{inv}(\ell^+ \ell^-)$
- $m_t \leftarrow M_{inv}(\ell^\pm J/\psi)$
- $m_t \leftarrow \sigma(pp \rightarrow t\bar{t}) = f(m_t)$



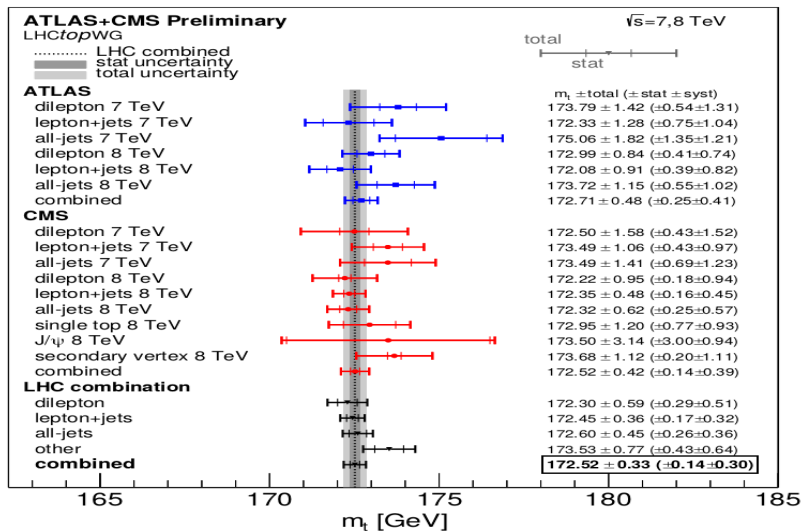
t-quark mass

CMS: $\sqrt{s} = 13$ TeV, five observables: $m_W, m_{\ell b}, m_t, m_{\ell b}/m_t, R_{\ell b}$,
ATLAS+CMS (15 measurements), $\sqrt{s} = 7, 8$ TeV

$$\begin{aligned}
 m_{t\text{op}}(\text{CMS}) &= 171.37 \pm 0.37 \text{ GeV, arXiv:2302.01967 (2023)} \\
 m_{t\text{op}}(\text{ATLAS+CMS}) &= 172.52 \pm 0.33 \text{ GeV, ATLAS-CONF-2023-066} \\
 \text{CMS} \Rightarrow \Delta m_t &= m_t - m_{\bar{t}} = -0.15 \pm 0.19(\text{stat}) \pm 0.09(\text{syst})
 \end{aligned}$$

RPP: *R.L. Workman et al. (Particle Data Group),
Prog.Theor.Exp.Phys.2022, 083C01 (2022) and 2023 update*

$$\begin{aligned}
 m_t(\text{RPP}) &= 172.69 \pm 0.30 \text{ GeV} && \text{direct measurements} \\
 m_t(\text{RPP}) &= 172.5 \pm 0.7 \text{ GeV} && \text{from } \sigma_{t\bar{t}}
 \end{aligned}$$

t-quark mass

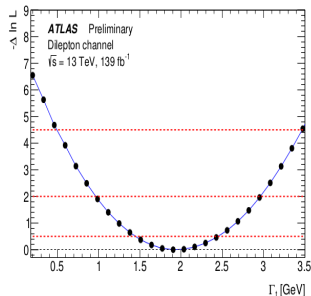
Measurements of the total decay width Γ_t

- conventional methods (peak value in $M_{inv}(bW)$) do not provide acceptable accuracy

- ATLAS - $M_{inv}^{exp}(j_b \ell^\pm)$ comparison with modeling events for different Γ_t
 $\Gamma_t = 0.1, 0.2, 0.3, \dots 5.0$ GeV, $\Delta\Gamma = 0.1$ GeV
 ATLAS Collaboration, ATLAS-CONF-2019-038

- CMS - indirect measurement
 CMS Collaboration, *Phys.Lett. B* 736 (2014), 33;
arXiv:1404.2292 [hep-ex]

$$\Gamma_t = \frac{\sigma_{t\text{-channel}}^{exp}}{\mathcal{B}(t \rightarrow Wb)} \times \frac{\Gamma^{th}(t \rightarrow bW)}{\sigma_{t\text{-ch}}^{th}}$$



| collaboration | Γ_t^{exp} , GeV | Γ_t^{SM} |
|---------------|---|-----------------|
| ATLAS | $\Gamma_t = 1.94^{+0.52}_{-0.49}$ | 1.39 GeV |
| CMS | $\Gamma_t = 1.36 \pm 0.02(stat.)^{+0.14}_{-0.11}$ (syst.) | |

Measurements of t -quark parameters

Measurements of events with $t\bar{t}$ -pair production with subsequent decays into different final states:

dileptons: $ee, e\mu, \mu\mu, \ell + jets$: $e + jets, \mu + jets$ and $\ell + \tau$ allow ones to measure decay probabilities through various channels (Br, %)

| channel | Br, % | SM | W^\pm (LEP) |
|---------------------------|--|-------------------|------------------|
| $t \rightarrow bj\bar{j}$ | 66.5 ± 0.4 (stat) ± 1.3 (syst) | 67.51 ± 0.007 | 67.48 ± 0.28 |
| $t \rightarrow be\nu$ | 13.3 ± 0.4 (stat) ± 0.5 (syst) | 12.72 ± 0.01 | 67.48 ± 0.20 |
| $t \rightarrow b\mu\nu$ | 13.4 ± 0.3 (stat) ± 0.5 (syst) | 12.72 ± 0.01 | 12.60 ± 0.18 |
| $t \rightarrow b\tau\nu$ | 7.0 ± 0.3 (stat) ± 0.5 (syst) | 7.05 ± 0.01 | 7.2 ± 0.12 |

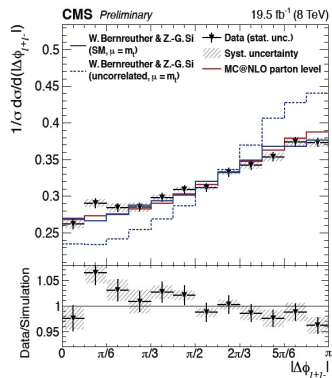
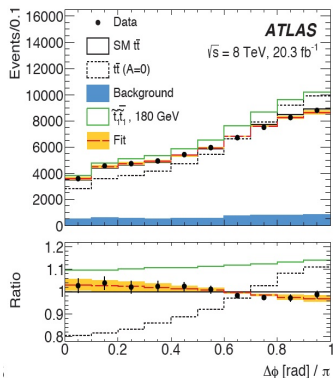
- CMS

$$pp \rightarrow t\bar{t}X, \quad t(\bar{t}) \rightarrow \ell^\pm j_b$$

$$R = \frac{\mathcal{B}(t \rightarrow bW)}{\sum_{q=s,d,s,b} \mathcal{B}(t \rightarrow qW)} = 1.014 \pm 0.003(\text{stat}) \pm 0.032(\text{syst})$$

Measurements of t -quark parameters

- polarized t -quark transmits the information about the spin to the decay products (angular distributions)
- “at the threshold” pair of $t\bar{t}$ quarks is produced mainly with identical helicities, and for large $\sqrt{S_{t\bar{t}}}$ - with opposite ones
- $pp \rightarrow t\bar{t}X$, $t \rightarrow l^+X$, $\bar{t} \rightarrow l^-X$: $\Delta\phi_{l^+l^-}$ sensitive to the presence of quark polarization

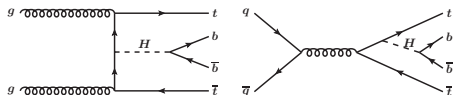


Top-quark and Higgs boson

- all Higgs boson couplings are evaluated within SM ($g, v, m_t, m_{W/Z}$)

$$f\bar{f}H \Leftrightarrow y_f = \sqrt{2} \frac{m_f}{v}; \quad y_t(m_t = 172.5 \text{ GeV}) = 0.99, \quad y_b(m_b = 4.5 \text{ GeV}) = 0.02$$

- $t\bar{t}$ -pair and Higgs production provides the direct measurement of y_t



- Higgs decay channels

| channel | $H \rightarrow b\bar{b}$ | $H \rightarrow WW/ZZ$ | $H \rightarrow gg$ | $H \rightarrow \gamma\gamma$ |
|---------|--------------------------|-----------------------|--------------------|------------------------------|
| Br | $\sim 58 \%$ | $\sim 24 \%$ | $\sim 8 \%$ | $\sim 0.2 \%$ |

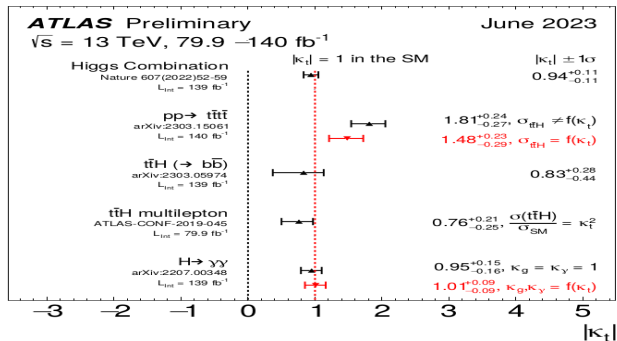
the main contribution to the Higgs inclusive cross section production comes from gluon annihilation $gg \rightarrow H$ (loop contribution)

Higgs boson and t -quark

y_t parameter can be measured both in pair and single production of t -quarks

$$pp \rightarrow t\bar{t}HX, pp \rightarrow tHX, H \rightarrow b\bar{b}, WW^*, ZZ^*, \tau^+\tau^-, \gamma\gamma$$

$$\kappa_t = y_t^{mes}/y_t^{SM} \approx 1.01 \pm 0.1$$



- there were obtained the limits on on the anomalous interaction of the Higgs boson with the t -quark

$$\mathcal{L} = -\frac{y_t}{\sqrt{2}} \bar{\psi}_t (C_\alpha \kappa_{Htt} + i s_\alpha \kappa_{Att} \gamma_5) \psi_t X_0$$

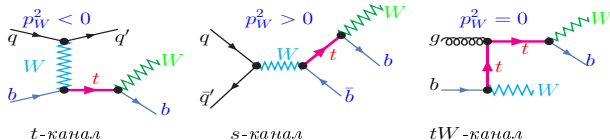
Electroweak t -quark production - "single" top

- The cross section production up to 50% to QCD

There are three processes of the single top production

$$\sigma_{t\text{-chan}}(t + \bar{t}) \simeq 218 \text{ pb}, \quad \sigma_{tW\text{-chan}}(t + \bar{t}) \simeq 70 \text{ pb}, \quad \sigma_{s\text{-chan}}(t + \bar{t}) \simeq 11 \text{ pb}$$

- production mechanisms depend on W -boson virtuality



- cross sections are calculated with NNLO accuracy, in this case, NNLO makes a small correction (mutual cancellation of QCD and EW)

- calculations can be done in 4F (without taking into account the contribution of the initial b -quarks) and 5F

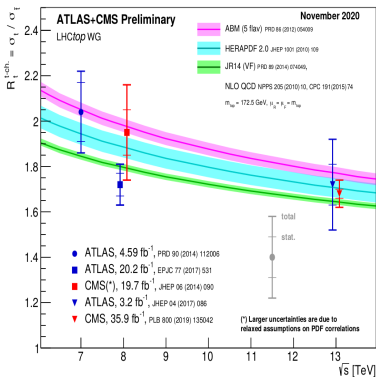
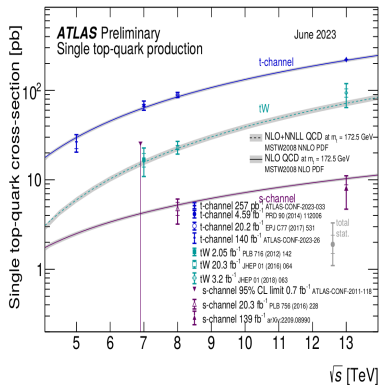


$$\sqrt{s} = 13 \text{ TeV}, \quad \sigma_{t\text{-chan}}(5Fs) = 134.0_{-0.6}^{+0.7} \text{ pb}, \quad \sigma_{t\text{-chan}}(4Fs) = 127.0_{-0.6}^{+2.0} \text{ pb}$$

Electroweak t -quark production

s -channel

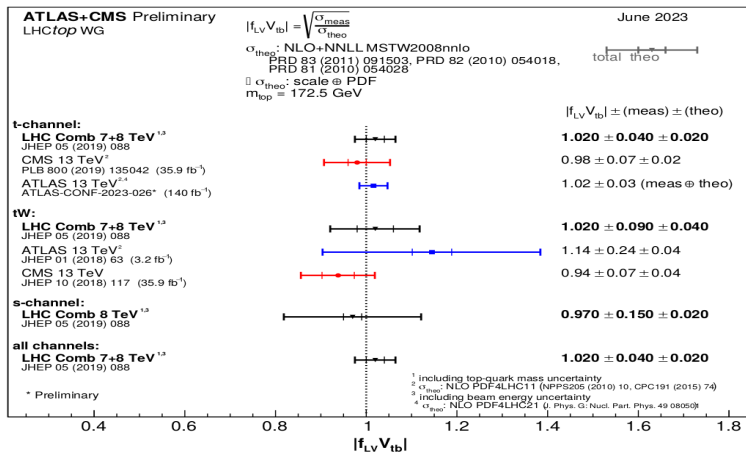
| | | |
|-------------------|--|---|
| <i>ATLAS</i> | $\sigma = 8.2 \pm 0.6(stat)_{-2.8}^{+3.4}(syst) = 8.2_{-2.9}^{+3.5}$ pb, | $\sigma^{theor} = 10.32_{-0.36}^{+0.40}$ pb |
| <i>theory(tW)</i> | $\sigma = 71.7 \pm 1.8(scale) \pm 3.4(PDF)$ pb | |
| <i>ATLAS(tW)</i> | $\sigma = 94 \pm 109(stat)_{-22}^{+28}(syst) \pm 2(lumi)$ pb | |
| <i>CMS(tW)</i> | $\sigma = 79.2 \pm 0.9(stat)_{-8.0}^{+2.7}(syst) \pm 1.2(lumi)$ pb | |



$|V_{tb}|$ measurements

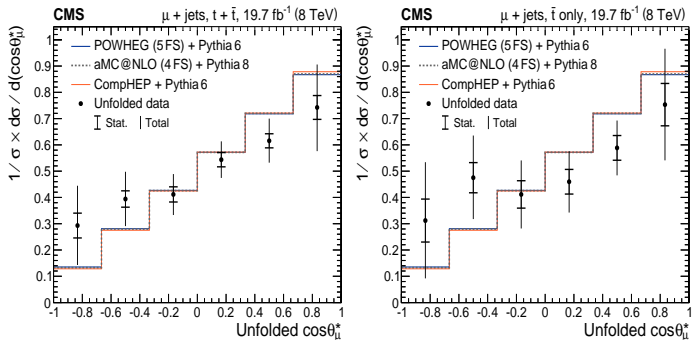
- cross-section of the electroweak t -quark production $\sigma_{EW}(tX) \propto |V_{tb}|^2$

$$\Rightarrow |V_{tb}|_{EW \text{ top}} = 1.019 \pm 0.028$$



Measuring t -quark polarization

- CMS: t -quark polarization in processes single top production: $bu \rightarrow d t (\rightarrow \ell^+ X)$
- distributions over $\cos \vartheta_{\mu}^*$ - between ℓ^{\pm} from decay t -quark and light quark
- correlations of the spin states in the production and decay of the t -quark

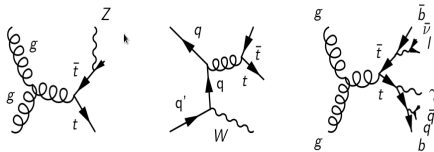


Top-quarks and W/Z -bosons associated production

vector bosons and t -quark associated production processes:

$$pp \rightarrow t\bar{t} + Z, t\bar{t} + W^\pm, t\bar{t} + \gamma, \quad \text{CDF} : \sigma(pp \rightarrow t\bar{t}\gamma)_{\sqrt{s}=2 \text{ TeV}} = 180 \pm 80 \text{ fb}$$

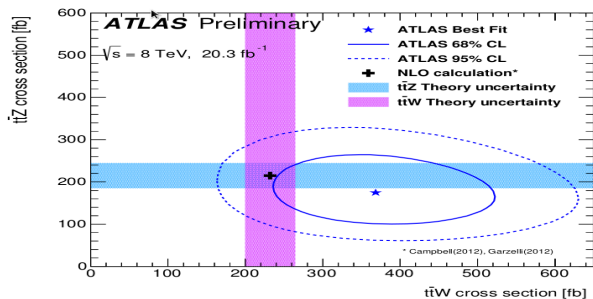
are sensitive to physics beyond SM



| $\sigma(t\bar{t}V)$ | ATLAS (fb) | CMS (fb) |
|--|---|---|
| $\sigma_{t\bar{t}\gamma}(7 \text{ TeV})$ | $2000 \pm 500(\text{stat}) \pm 700(\text{syst})$ | |
| $\sigma_{t\bar{t}\gamma}(8 \text{ TeV})$ | | $2400 \pm 200(\text{stat}) \pm 600(\text{syst})$ |
| $\sigma_{t\bar{t}Z/\gamma^*}(7 \text{ TeV})$ | < 700 | $280^{+140}_{-110}(\text{stat})^{+60}_{-30}(\text{syst})$ |
| $\sigma_{t\bar{t}Z/\gamma^*}(8 \text{ TeV})$ | $150^{+55}_{-50}(\text{stat}) \pm 21(\text{syst})$ | $200 \pm 90(\text{total})$ |
| $\sigma_{t\bar{t}W}(8 \text{ TeV})$ | $300^{+120}_{-100}(\text{stat})^{+70}_{-40}(\text{syst})$ | $170^{+110}_{-100}(\text{total})$ |

Top-quarks and W/Z -bosons associated production

| $\sqrt{s} = 8 \text{ TeV}$ | $\sigma(\text{SM}), \text{ fb}$ | ATLAS (fb) | CMS (fb) |
|----------------------------|---------------------------------|--|---------------------|
| $t\bar{t} + \gamma$ | 1880 ± 500 | | |
| $t\bar{t}Z / \gamma^*$ | 215 ± 30 | $176_{-48}^{+52}(\text{stat}) \pm 24(\text{syst})$ | 242_{-55}^{+65} |
| $t\bar{t}W$ | 232 ± 32 | $369_{-79}^{+86}(\text{stat}) \pm 44(\text{syst})$ | 382_{-102}^{+117} |



Search for the New physics beyond the Standard Model

In the top quark sector, New physics may manifest itself in the following

- ◇ rare (within the SM) decays of t -quarks;
- ◇ deviations in t -quark production cross sections (within the SM framework);
- ◇ production of t -quarks due to very rare reactions in the SM;
- ◇ decays of t -quarks through channels absent in the SM;
- ◇ new particles decaying into final states containing t and/or \bar{t} -quark and possibly other particles and resonances

- Numerous SM extensions \implies various predictions in the t -quark sector with their a specific set of interaction types and parameters (coupling constants, masses of new objects)
- different scenarios \implies processes with identical final states

- **Effective field theory formalism** - effective (phenomenological) Lagrangian \mathcal{L}_{EFT} gauge-invariant with respect to to calibration group SM

$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \kappa_4 \bar{\psi}_q \hat{O}^{(4)} \psi_t + \frac{\kappa_6}{\Lambda^2} \bar{\psi}_q \hat{O}^{(6)} \psi_t + \dots$$

- Experimental results are presented in the form restrictions
 - ◇ κ/Λ - values of anomalous constants interaction
 - ◇ in the form of limits on probability rare decays of the t -quark

Anomalous $gt\bar{t}$ and tWb interactions

- $gt\bar{t}$: deviations from the SM can manifest themselves in energy and angle distributions phenomenological Lagrangian (with anomalous chromomagnetic moment):

$$\mathcal{L}_{eff} = \mathcal{L}_{SM} - \frac{\tilde{\mu}_t}{2} \bar{\psi}_t \sigma^{\mu\nu} \psi_t G_{\mu\nu}^a \quad \text{experiment} \Rightarrow -0.50 < \text{Re}(\tilde{\mu}_t) < 0.070 (95\% CL)$$

- tWb : effective interaction Lagrangian

$$\mathcal{L}_{EFT} = \frac{g}{\sqrt{2}} \bar{b} \gamma^\mu (f_V^L P_L + f_V^R P_R) t W_\mu^- + \frac{g}{\sqrt{2}} \bar{b} \frac{\sigma^{\mu\nu}}{2M_W} (f_T^L P_L + f_T^R P_R) t W_{\mu\nu}^- + h.c.$$

where $W_{\mu\nu}^- = \partial_\mu W_\nu^- - \partial_\nu W_\mu^-$, within SM one has: $f_V^L = V_{tb}$; $f_V^R = f_T^L = f_T^R = 0$

$$|f_V^R| < 0.16, \quad |f_T^L| < 0.057, \quad |f_T^R| < 0.048 \quad \text{at } 95\% \text{ CL}$$

Flavor changing neutral currents – FCNC: tVq

FCNC interactions tVq , $V = g, \gamma, Z, H$ strongly suppressed within SM

| | SM | two-Higgs | SUSY | "exotic" quarks |
|-----------------------------|---------------------|----------------|----------------|-------------------------|
| $BR(t \rightarrow qg)$ | 5×10^{-11} | $\sim 10^{-5}$ | $\sim 10^{-3}$ | $\sim 5 \times 10^{-4}$ |
| $BR(t \rightarrow q\gamma)$ | 5×10^{-13} | $\sim 10^{-7}$ | $\sim 10^{-5}$ | $\sim 10^{-5}$ |
| $BR(t \rightarrow qZ)$ | $\sim 10^{-13}$ | $\sim 10^{-6}$ | $\sim 10^{-4}$ | $5 \times \sim 10^{-2}$ |

| | SM | SUSY | MSSM | 2HDM |
|------------------------|---------------------|-----------|--------------------|-----------|
| $BR(t \rightarrow Hc)$ | 3×10^{-15} | 10^{-6} | 10^{-5} | 10^{-3} |
| $BR(t \rightarrow Hu)$ | 2×10^{-17} | 10^{-6} | 8×10^{-5} | 10^{-4} |

model independent analysis. The phenomenological Lagrangian

$$\begin{aligned}
 \mathcal{L}_{FCNC} = & -e \sum_{q=u,c} \frac{\kappa_q^\gamma}{\Lambda} \bar{t} \sigma^{\mu\nu} (f_q^\gamma + ih_q^\gamma \gamma_5) q A_{\mu\nu} - g_s \sum_{q=u,c} \frac{\kappa_q^g}{\Lambda} \bar{t} \sigma^{\mu\nu} t^a (f_q^g + ih_q^g \gamma_5) q G_{\mu\nu}^a \\
 & - \frac{g}{2 \cos \theta_W} \sum_{q=u,c} \kappa_q^Z \bar{t} \gamma^\mu (f_q^Z - h_q^Z \gamma_5) q Z_\mu - \frac{g}{2 \cos \theta_W} \sum_{q=u,c} \frac{\tilde{\kappa}_q^Z}{\Lambda} \bar{t} \sigma^{\mu\nu} (\tilde{f}_q^Z + i\tilde{h}_q^Z \gamma_5) q Z_{\mu\nu}
 \end{aligned}$$

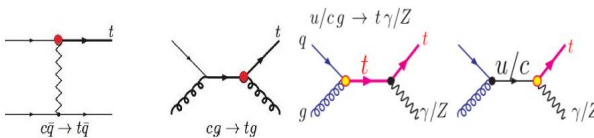
FCNC tq , $t\gamma q$, tZq , tHq

- two scenarios for searching (FCNC/SM $\ll 1$)
- ◇ $t\bar{t}$ -pair production with the subsequent FCNC decay

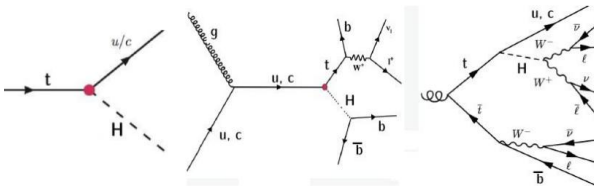
$$pp \rightarrow t\bar{t} : t \rightarrow qg, t \rightarrow q\gamma, t \rightarrow qZ$$

- ◇ t -quark production due to FCNC with subsequent SM decay

$$g u(c) \rightarrow t; uu \rightarrow tt; cg \rightarrow tg; qg \rightarrow t\gamma/Z, \dots t \rightarrow bW$$



- ◇ tHq : $t \rightarrow H u/c \rightarrow$ Higgs + FCNC



Summary for FCNC $t \rightarrow q H/g/\gamma/Z$

the best constraints on anomalous FCNC interactions in the t -quark sector

| channel | \mathcal{B} (95% CL) | | mode | \mathcal{B} (95% CL) | |
|--------------------------|------------------------|-----------------------|--------------------------|------------------------|-----------------------|
| | ATLAS | CMS | | ATLAS | CMS |
| $t \rightarrow Hu$ | 6.9×10^{-4} | 1.9×10^{-4} | $t \rightarrow Hc$ | 9.4×10^{-4} | 7.3×10^{-4} |
| $t \rightarrow gu$ | 4.0×10^{-5} | 2.0×10^{-5} | $t \rightarrow gc$ | 20.0×10^{-5} | 41×10^{-5} |
| $t \rightarrow \gamma u$ | | 0.95×10^{-5} | $t \rightarrow \gamma c$ | | 1.51×10^{-5} |
| $t \rightarrow Zu$ | 6.2×10^{-5} | 22×10^{-5} | $t \rightarrow Zc$ | 12×10^{-5} | 44×10^{-5} |

Lepton-flavor violation (CMS arXiv:22060159)

$t \rightarrow e^\pm \mu^\mp q$ due to vector (V), scalar (S) tensor (T) interactions

| channel | \mathcal{B} (95% CL) | | |
|---------------------------------|------------------------|----------------------|-----------------------|
| | V | S | T |
| $t \rightarrow e^\pm \mu^\mp u$ | 1.3×10^{-7} | 0.7×10^{-7} | 2.5×10^{-7} |
| $t \rightarrow e^\pm \mu^\mp c$ | 13.1×10^{-7} | 8.9×10^{-7} | 25.9×10^{-7} |

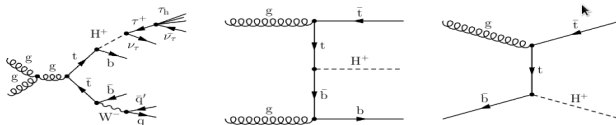
Charged Higgs H^\pm and rare processes with t -quarks

- interaction Lagrangian of the charged Higgs H^\pm (MSSM)

$$\mathcal{L} = \frac{g}{\sqrt{2}M_W} H^+ \{ V_{ud} \bar{u} (m_u \cot \beta P_L + m_d \tan \beta P_R) d + \bar{\nu} (\tan \beta m_\ell P_R) \ell \}, \quad P_{L/R} = 1/2 (1 \mp \gamma^5)$$

- two regions of the charged Higgs mass

$$m_{H^\pm} = 80 - 160 \text{ GeV} : t \rightarrow H^\pm b \quad m_{H^\pm} > 180 \text{ GeV} : pp \rightarrow \bar{t} H^\pm b$$



- charged Higgs is excluded in the regions:

$$\tan \beta < \mathcal{O}(1)$$

$$\tan \beta \gtrsim 1$$

$$m_H < 180 \text{ GeV} \text{ and } m_H > 180 \text{ GeV}$$

$$m_H (90 \div 2000) \text{ GeV}$$

Search for heavy particles decaying into t -quarks

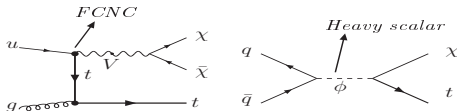
- CMS: arXiv:2310.19893

new heavy charged vector boson $W' \rightarrow tb$, $\sqrt{s} = 13$ TeV, $\mathcal{L}_{int} = 138 \text{ fb}^{-1}$

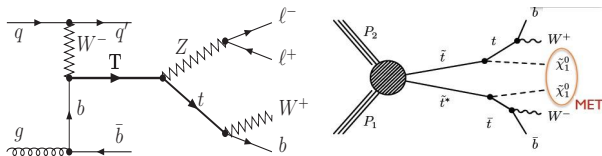
Multiple hypotheses are considered for the new particle mass, width, and chirality

| $\Gamma(W')/M(W')$ | right-handed | left-handed |
|--------------------|--------------|-------------|
| 1% | < 4.3 TeV | < 3.9 TeV |
| 10% | < 2.7 TeV | < 2.5 TeV |

- SUSY predicts a large number of new reactions with t quarks with presence of large "lost" energy carried away by the new neutral particle. For example, [single \$t\$ -quark](#) production with a large "lost" energy ("mono-top") \Rightarrow search for "dark matter"

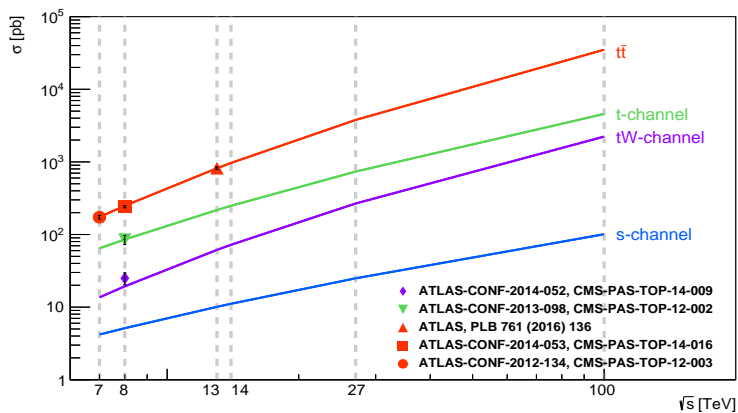


- "vector" T -quarks: $T \rightarrow tZ$. SUSY \Rightarrow ($\tilde{t} \rightarrow t\chi_1^0$, $\text{spin}(\tilde{t}) = 0$), constraint: $m_t \leq m_{stop} < 195$ GeV



Top-quark physics at future hadronic colliders

| collider | \sqrt{s} , TeV | \mathcal{L} , $\text{cm}^{-2} \cdot \text{c}^{-1}$ | $\int \mathcal{L}$, ab^{-1} | $\langle \mu \rangle$ |
|----------|------------------|--|---------------------------------------|-----------------------|
| LHC | 7-13 | $\approx 10^{34}$ | 0.3 | 10-40 |
| HL-LHC | 14 | 10^{35} | 3 | 140-200 |
| HE-LHC | 27 | 2.5×10^{35} | 12 | 800 |
| FCC-hh | 100 | 3×10^{35} | 30 | 500-1000 |



Summary of the top-quark properties (RPP)

R.L. Workman et al. (Particle Data Group), *Prog.Theor.Exp.Phys.* **2022**, 083C01 and 2023

$$\begin{aligned}
 m_t &= 172.69 \pm 0.30 \text{ GeV} && \text{direct measurements} \\
 m_t &= 172.5 \pm 0.7 \text{ GeV} && \text{from } \sigma_{t\bar{t}} \\
 m_t - m_{\bar{t}} &= 0.15 \pm 0.20 \text{ GeV} \\
 \Gamma_{tot} &= 1.42^{+0.19}_{-0.15} \text{ GeV} \\
 \Gamma(Wb)/\Gamma(Wq(q = b, s, d)) &= 0.957 \pm 0.034
 \end{aligned}$$

t -quark decay modes (Br = Γ_j/Γ_{tot} and CL - confidence level)

| decay mode | Br | CL |
|--|-----------------------------|-----|
| $e\nu_e b$ | $(11.10 \pm 0.30)\%$ | |
| $\mu\nu_\mu b$ | $(11.40 \pm 0.20)\%$ | |
| $\tau\nu_\tau b$ | $(10.7 \pm 0.5)\%$ | |
| $q\bar{q}b$ | $(66.5 \pm 1.4)\%$ | |
| $\gamma q(q = u, c)$ | $< 1.8 \times 10^{-4}$ | 95% |
| $Zq(q = u, c)$ | $< 5 \times 10^{-4}$ | 95% |
| $Hu(c)$ | $< 1.9(7.3) \times 10^{-4}$ | 95% |
| $\ell^+ q\bar{q}'(q = d, s, b; q' = u, c)$ | $< 1.6 \times 10^{-3}$ | 95% |
| $e^\pm \mu^\mp c$ | $< 8.9 \times 10^{-7}$ | |
| $e^\pm \mu^\mp u$ | $< 7 \times 10^{-8}$ | |

Thank you very much !